

Lect. 2: Light as EM Waves (Chap. 8 in Cheng)

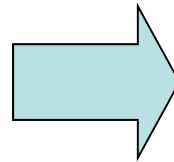
In the beginning, God said

$$\nabla \times \bar{\mathbf{E}} = -\frac{\partial \bar{\mathbf{B}}}{\partial t}$$

$$\nabla \times \bar{\mathbf{H}} = \bar{\mathbf{J}} + \frac{\partial \bar{\mathbf{D}}}{\partial t}$$

$$\begin{aligned} \nabla \cdot \bar{\mathbf{D}} &= \rho \\ \nabla \cdot \bar{\mathbf{B}} &= 0 \end{aligned} \quad \left\{ \begin{array}{l} \bar{\mathbf{D}} = \epsilon \bar{\mathbf{E}} \\ \bar{\mathbf{B}} = \mu \bar{\mathbf{H}} \end{array} \right.$$

Material Parameters $\left\{ \begin{array}{l} \epsilon : \textit{permittivity} \\ \mu : \textit{permeability} \end{array} \right.$



$$\nabla^2 \bar{\mathbf{E}} = \mu\epsilon \frac{\partial^2 \bar{\mathbf{E}}}{\partial t^2} \quad \nabla^2 \bar{\mathbf{H}} = \mu\epsilon \frac{\partial^2 \bar{\mathbf{H}}}{\partial t^2}$$

Wave Equations
(source free, uniform medium)

Then, there was light!

→ Governs light propagation

→ Maxwell's Equations

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Solutions for Wave Equations

$$\nabla^2 \bar{\mathbf{E}} = \mu\epsilon \frac{\partial^2 \bar{\mathbf{E}}}{\partial t^2}$$

Assume a plane-wave solution. For example, $\bar{\mathbf{E}} = \bar{x}E_0 e^{j(\omega t - kz)}$

$$\nabla^2 \bar{\mathbf{E}} = \bar{x}(-k^2)E_0 e^{j(\omega t - kz)}$$

$$k = \frac{2\pi}{\lambda} \quad \omega = 2\pi f$$

$$\mu\epsilon \frac{\partial^2 \bar{\mathbf{E}}}{\partial t^2} = \bar{x}\mu\epsilon(-\omega^2)E_0 e^{j(\omega t - kz)}$$

$$f \cdot \lambda = \frac{1}{\sqrt{\mu\epsilon}} \quad \text{Speed of light!}$$

$$k^2 = \mu\epsilon\omega^2$$

$$\therefore k = \omega\sqrt{\mu\epsilon}$$

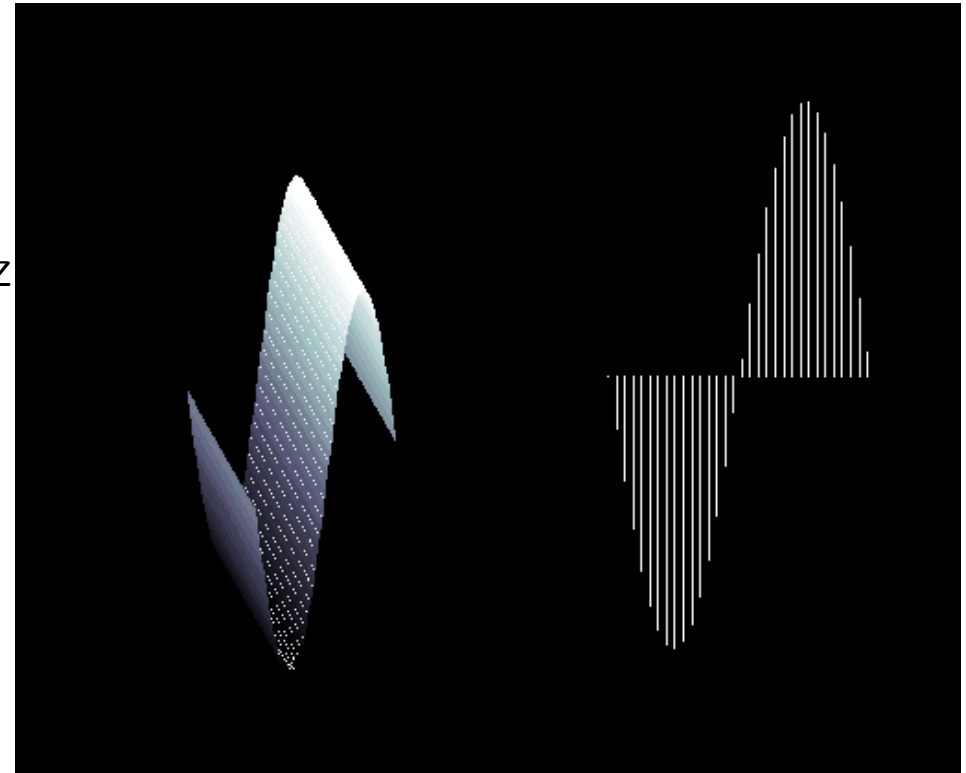
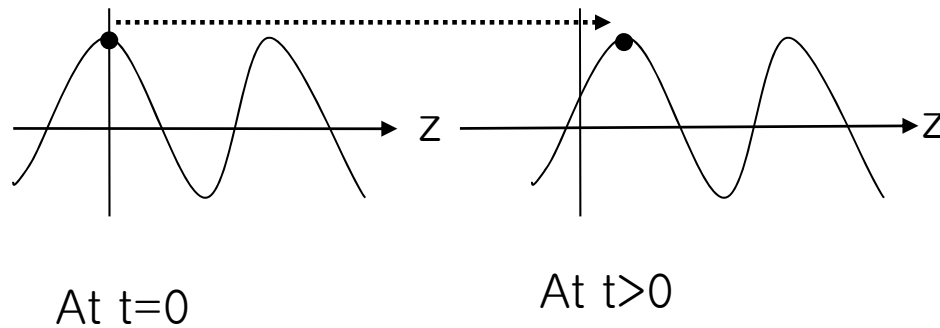
Why use plane-wave solutions?

(exponential solutions, time-harmonic solutions, phasor notation, ...)

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How does the plane-wave solution look like?

For physical representation, $\text{Re} \left[\bar{x}E_0 e^{j(\omega t - kz)} \right] = \bar{x}E_0 \cos(\omega t - kz)$



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How about H-field?

$$\nabla^2 \bar{\mathbf{E}} = \mu\epsilon \frac{\partial^2 \bar{\mathbf{E}}}{\partial t^2} \quad \bar{\mathbf{E}} = \bar{x} E_0 e^{j(\omega t - kz)}$$

It can be shown from Maxwell's Equations,

$$\bar{\mathbf{H}} = \bar{y} H_0 e^{j(\omega t - kz)}$$

Direction of E, H fields?

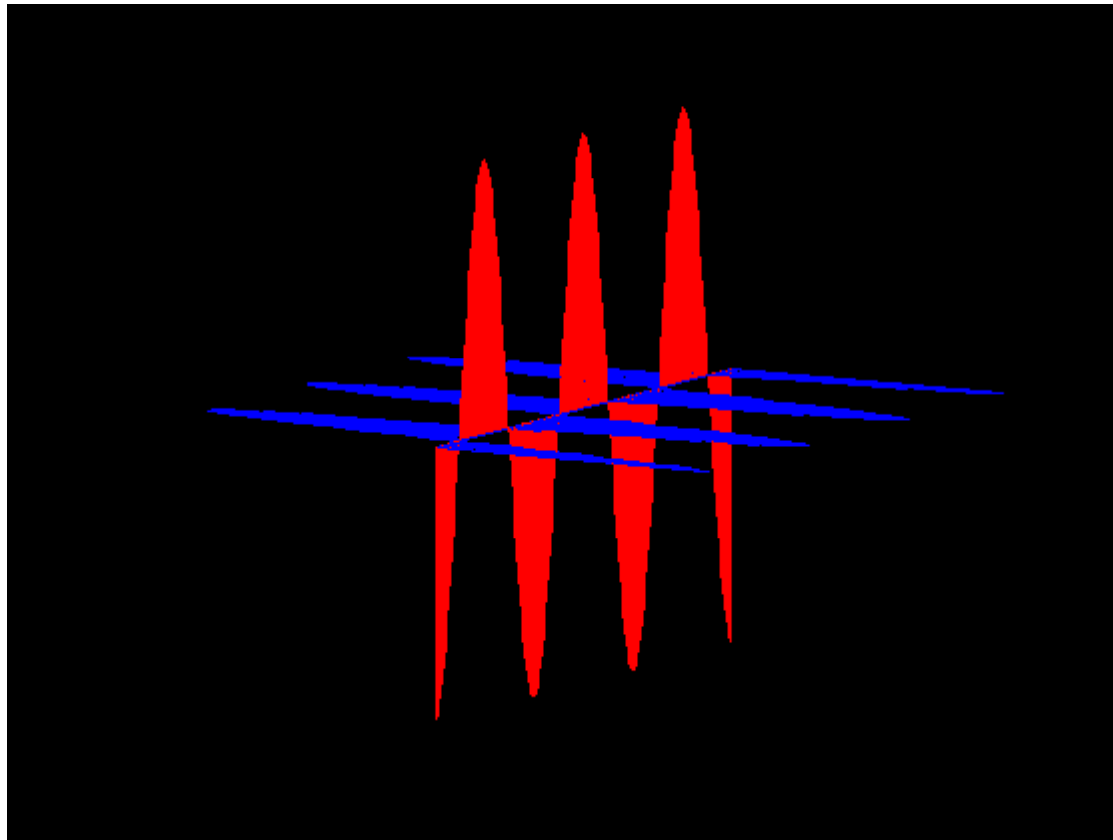
Direction of propagation?

Speed of propagation?

$$\frac{E_0}{H_0} = \sqrt{\frac{\mu}{\epsilon}} = \eta \text{ } [\Omega]$$

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How does the plane-wave solution look like?



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When a wave is propagating into

Plane wave solutions

+z direction: $e^{j(\omega t - kz)}$

-z direction: $e^{j(\omega t + kz)}$

+y direction: $e^{j(\omega t - ky)}$

Any direction? $e^{j\omega t} e^{-jk_x x} e^{-jk_y y} e^{-jk_z z} = e^{j(\omega t - \bar{k} \cdot \bar{R})}$

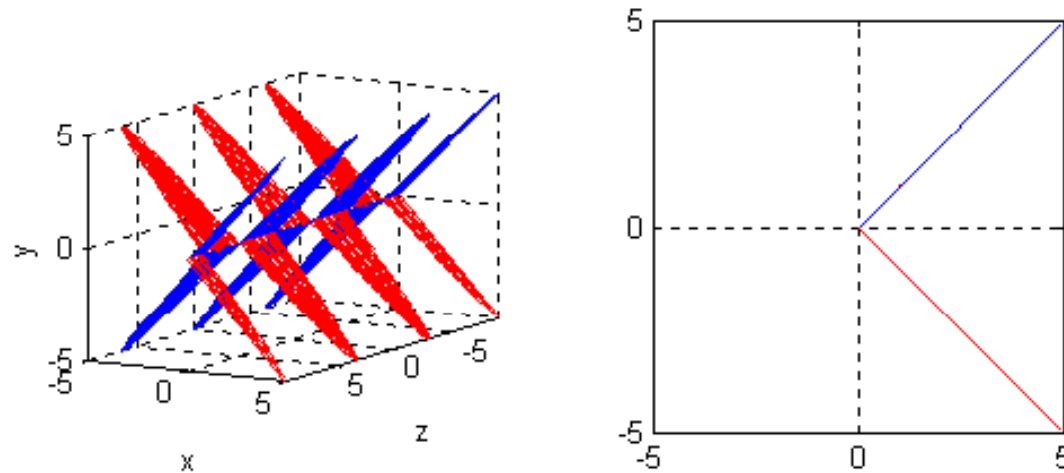
$$\bar{k} = \bar{x}k_x + \bar{y}k_y + \bar{z}k_z \quad \bar{R} = \bar{x}x + \bar{y}y + \bar{z}z$$

$$|\bar{k}| = \frac{2\pi}{\lambda}, \quad \angle \bar{k}: \text{direction of propagation}$$

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Polarization: Change of E-field direction with time

$$\text{Linear Polarization } \bar{E} = (\bar{x}E_0 + \bar{y}E_0)e^{j\omega t} e^{jkz}$$

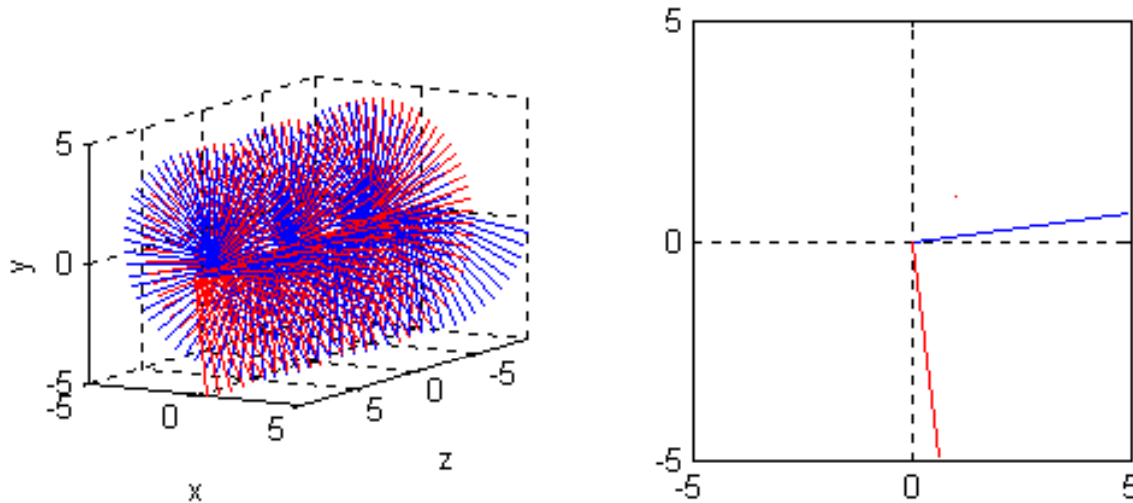


E_x and E_y in-phase

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Circular Polarization

$$\vec{E} = (\bar{x}E_0 + \bar{y}jE_0) e^{j\omega t} e^{jkz}$$

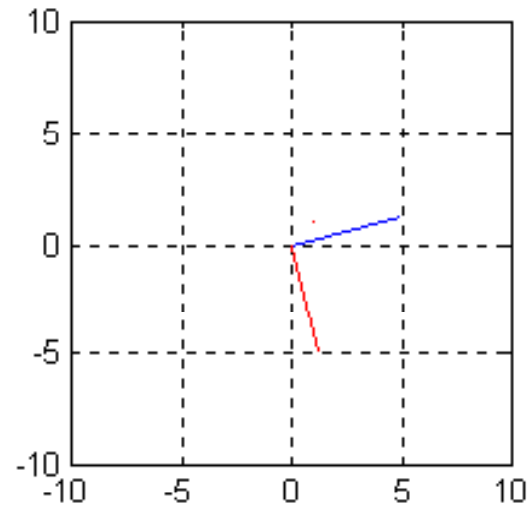
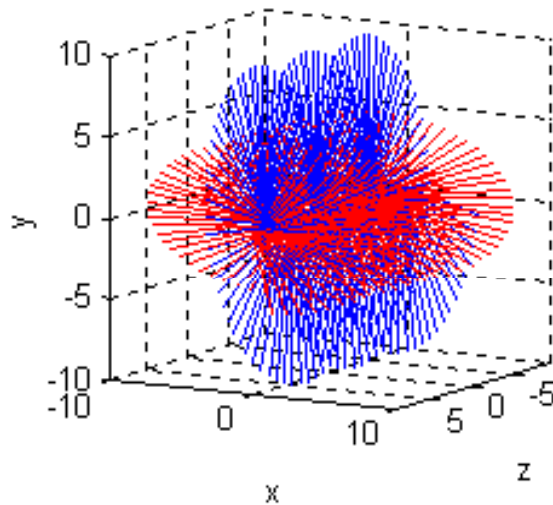


E_x and E_y out-of-phase Handedness?

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Elliptical Polarization

$$\bar{E} = (\bar{x}E_0 + \bar{y}j2E_0) e^{j\omega t} e^{jkz}$$



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Homework: Due on 9/13 before Tutorial

A uniform plane wave propagating in a dielectric medium has the E-field given as

$$\vec{E}(t, z) = \bar{x}2 \cos(10^8 t - z/\sqrt{3}) + \bar{y} \sin(10^8 t - z/\sqrt{3})$$

- (a) What is the dielectric constant of the medium?
- (b) What type of polarization does the wave have?
- (c) What is the corresponding H-field?